Breaking Wave Statistics with the multi-layer model

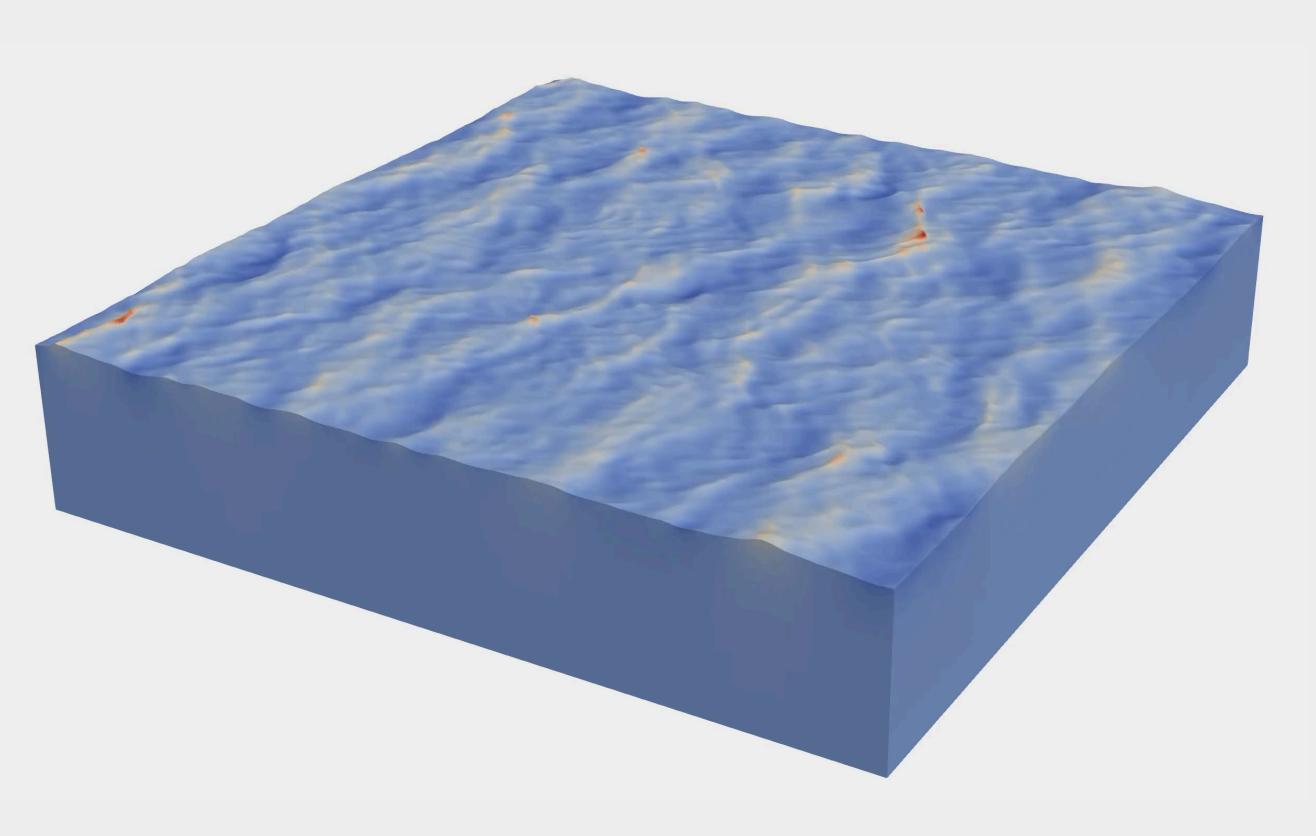
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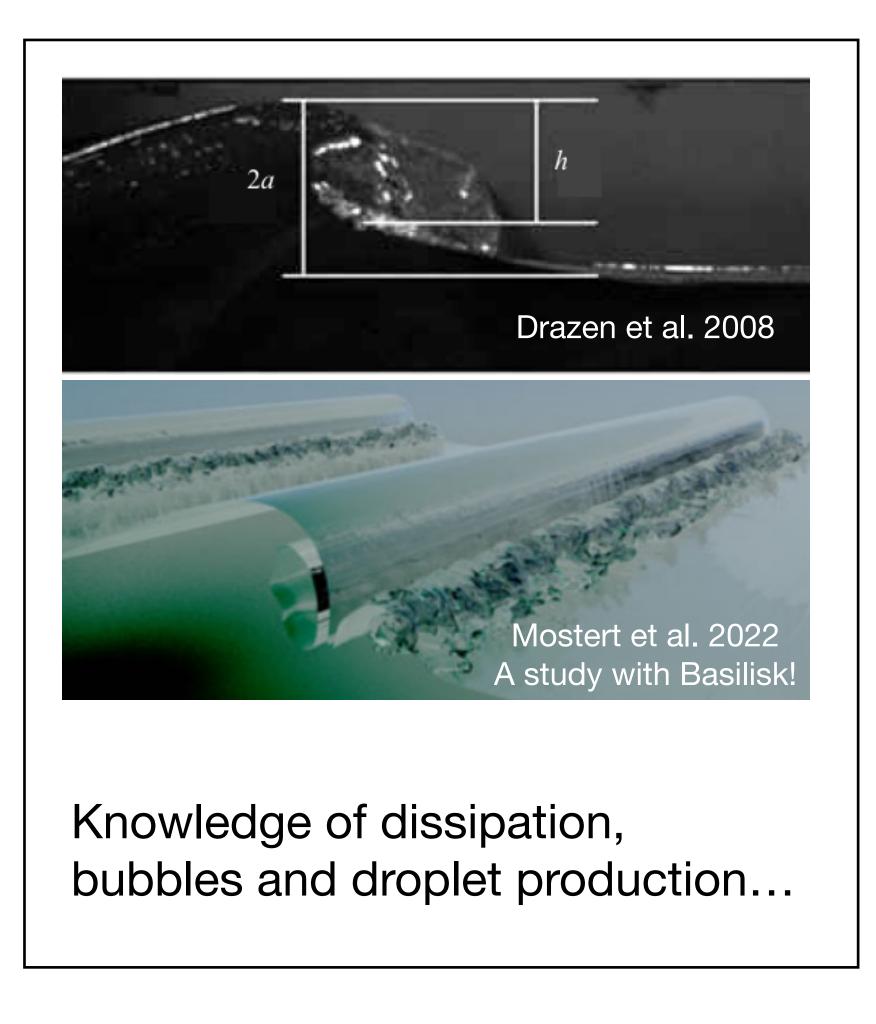
2023/07/05 Basilisk User Meeting







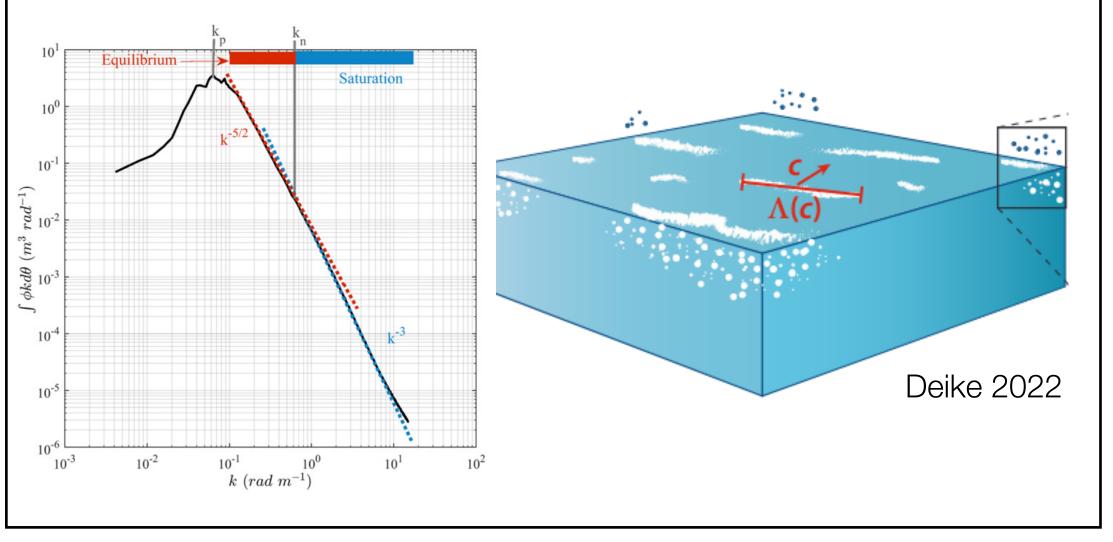
Wave breaking - from single breakers to broadband wave field



Goal: study an ensemble of breaking waves and their statistics. Breaking statistics: how frequently waves break and at what scales.

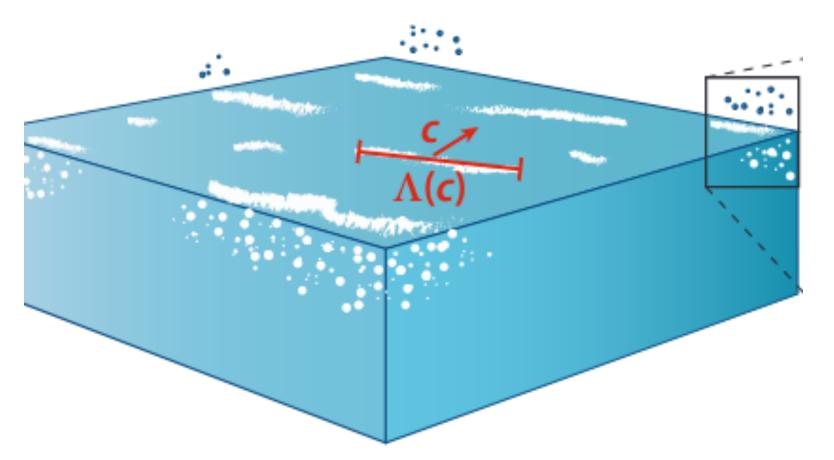


Wave breaking: enhanced mass transfer; contribute to upper ocean current and mixing.





Concept of breaking distribution



Total length of breaking fronts per unit surface area: $L = \int \Lambda(\boldsymbol{c}) d\boldsymbol{c}$

Larger c - smaller underlying k - stronger breaker

Fraction of total surface area turned over per unit time: $R = \int c \Lambda(\mathbf{c}) d\mathbf{c}$ Fractional whitecap coverage: $W \propto \int c^2 \Lambda(\mathbf{c}) d\mathbf{c}$ Rate of air entrainment per unit surface area: $V_a \propto \int c^3 \Lambda(\mathbf{c}) d\mathbf{c}$ Momentum flux per unit surface area: $M \propto \int c^4 \Lambda(\mathbf{c}) d\mathbf{c}$ Energy dissipation per unit surface area: $E \propto \int c^5 \Lambda(\mathbf{c}) d\mathbf{c}$

Romero 2019

Q: How do we predict breaking distribution $\Lambda(c)$ from a particular wave spectrum $\phi(k)$?

- Theories (Phillips 1985): for a typical wind wave spectrum $\Lambda(c) \propto u_*^3 g c^{-6}$
- Limited field campaigns (Gemmrich 2008, Kleiss & Melville 2010, Sutherland & Melville 2013, Deike 2017, etc.) proposed empirical corrections
- The breaking distribution and the associated dissipation in operational wave models are either **empirical** or **heavily tuned**
- **Numerical challenges**: highly non-linear, multi-scale, intermittent in time and space (need to integrate for a relatively long time).

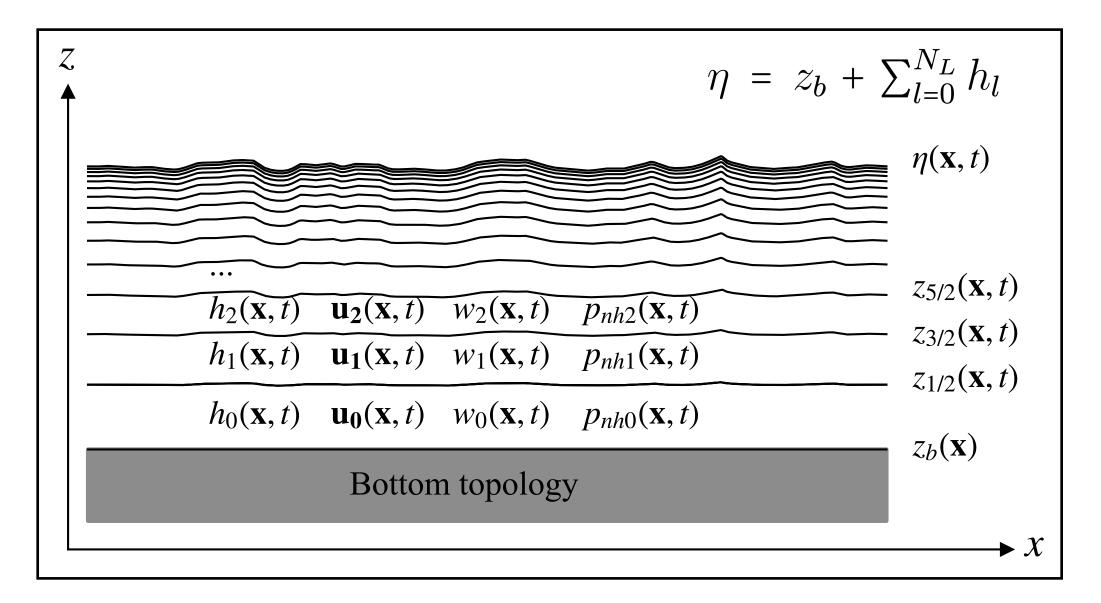
Our approach: p<u>hase-resolved</u> broadband wave field simulations in the <u>physical space</u> that permit <u>strong non-</u> <u>linearity</u> (with reasonable computational cost.)



The multi-layer model

- Multi-layer model <u>http://basilisk.fr/src/layered/README</u> (Popinet 2020);
- Discretization: horizontally Eulerian and vertically Lagrangian, with generalized vertical coordinate.

Layer structure illustration



of grid points: 1024*1024*15

(Depth-integrated) NS equation

$$\frac{\partial h_l}{\partial t} + \nabla_H \cdot (h\boldsymbol{u})_l = 0$$

$$\frac{\partial (h\boldsymbol{u})_l}{\partial t} + \nabla_H \cdot (h\boldsymbol{u}\boldsymbol{u})_l = -gh_l \nabla_H \eta - \nabla_H (h\phi)_l + [\phi \nabla_H z]_l$$

$$+ [\nu_1 \partial_z \boldsymbol{u}]_l + \nu_2 \nabla_H^2 \boldsymbol{u}$$

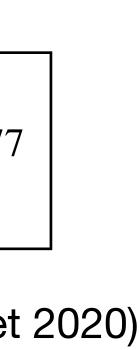
$$\frac{\partial (hw)_l}{\partial t} + \nabla_H \cdot (hw\boldsymbol{u})_l = -[\phi]_l + [\nu_1 \partial_z w]_l + \nu_2 \nabla_H^2 w$$

$$\nabla_H \cdot (h\boldsymbol{u})_l + [w - \boldsymbol{u} \cdot \nabla_H z]_l = 0 \qquad \phi: \text{ non-hydrostatic pressure}$$

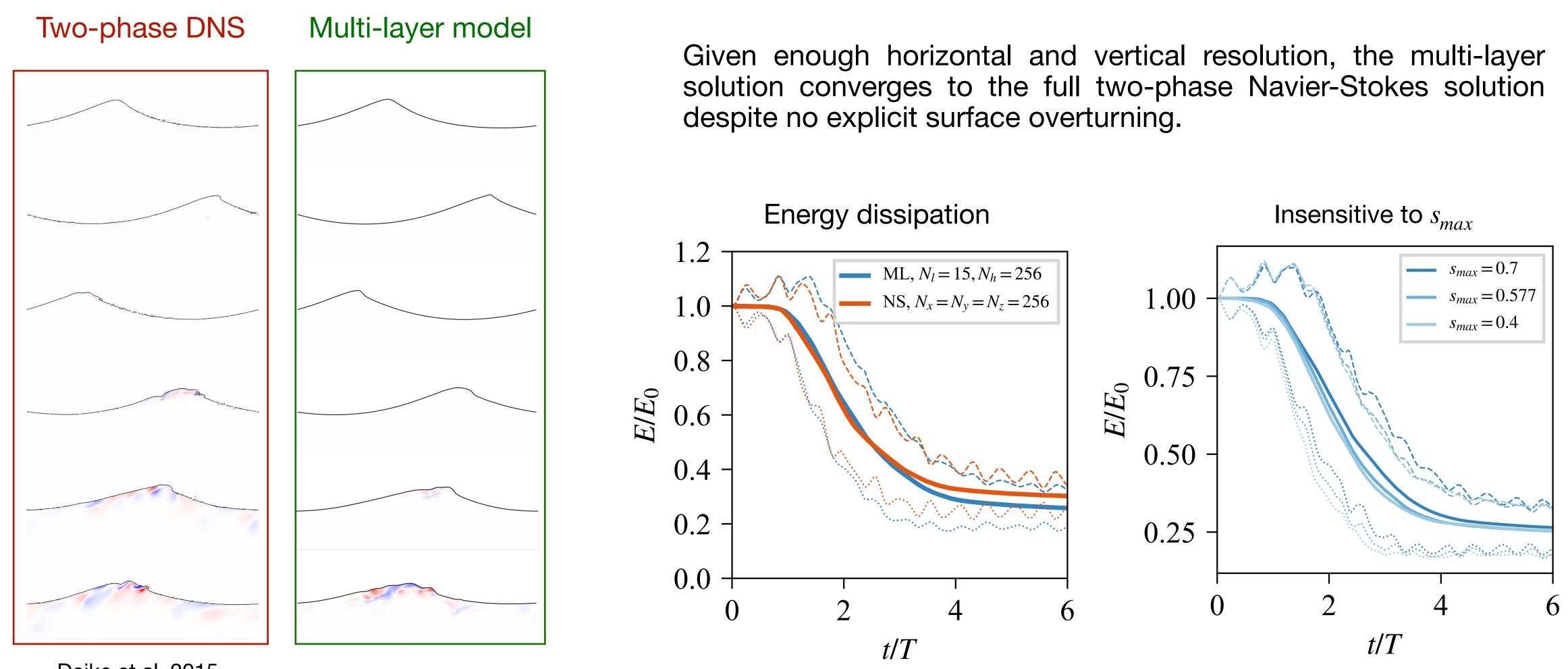
Combined with a gradient-limiter for breaking

$$\frac{\partial z}{\partial x} = \begin{cases} \frac{\partial z}{\partial x}, & |\frac{\partial z}{\partial x}| \leq s_{\max} \\ \operatorname{sign}(\frac{\partial z}{\partial x})s_{\max}, & |\frac{\partial z}{\partial x}| > s_{\max}. \end{cases} \quad s_{max} = 0.57$$

Numerical scheme: advection - projection - remapping (Popinet 2020)



Multil-layer v.s. NS for a single breaker



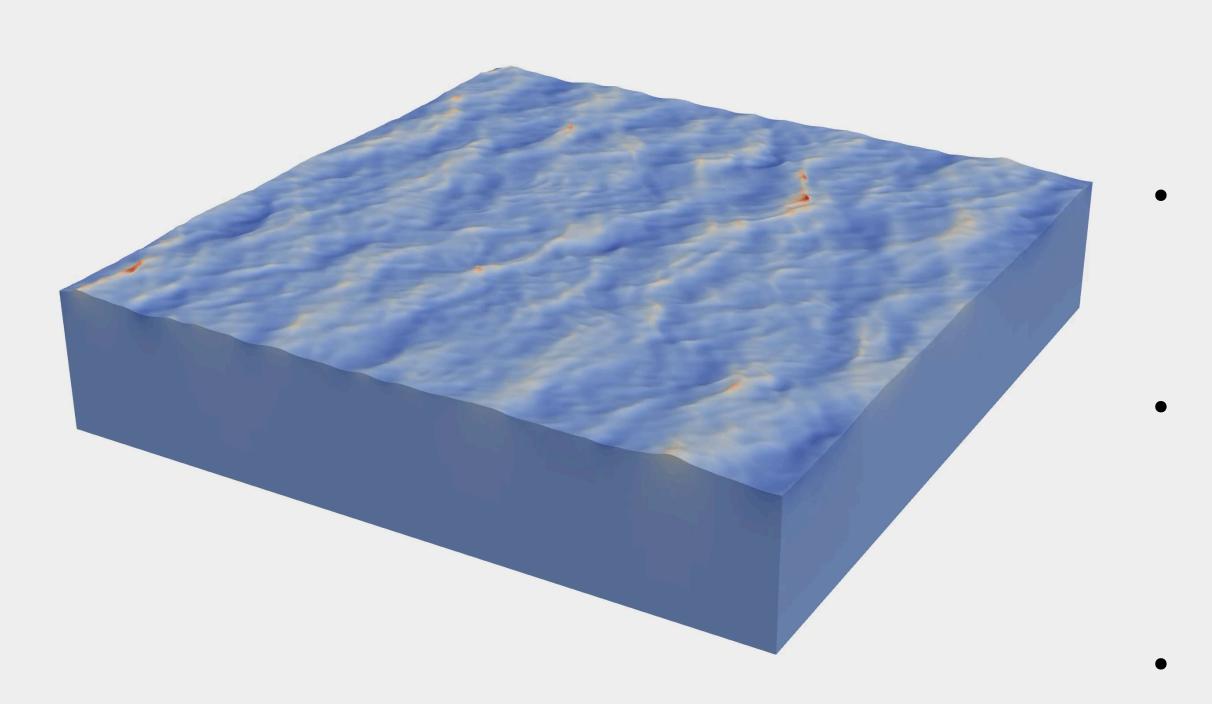
Deike et al. 2015, Mostert et al. 2022, etc.

About 40 times faster.

Simulation configuration

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Choose a typical wind wave energy spectrum:

$$\phi(k) = Pg^{-1/2}k^{-2.5} \exp[-1.25(k_p/k)^2]$$

Directional spreading:

$$F(k,\theta) = \left(\frac{\phi(k)}{k}\right) \cos^{N}(\theta) / \int_{-\pi/2}^{\pi/2} \cos^{N}(\theta) d\theta$$

Create one random realization of the wave spectrum: superposition of linear waves with random phase as initial conditions:

$$\eta = \sum_{i,j} a_{ij} \cos(\psi_{ij}), \quad a_{ij} = [2F(k_{xi}, k_{yj})dk_x dk_y]^{1/2}, \quad \psi_{ij} = k_x x + k_y y + \psi_{ra}$$

Breaking modeled with a gradient-limiter (no surface overturning). No wind forcing.

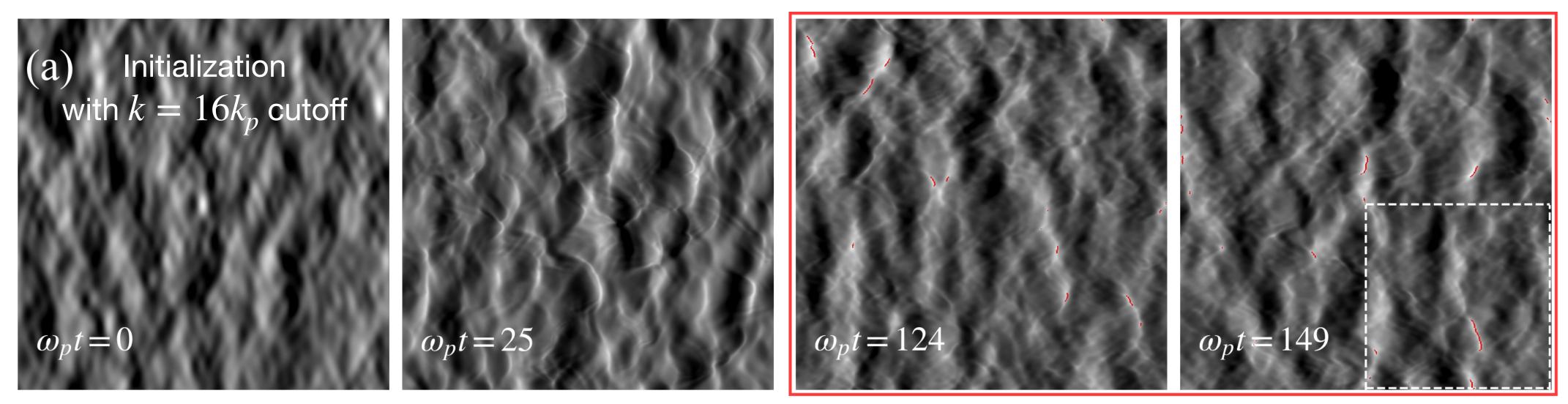
8 hrs on 256 cores for \sim 40 peak wave period.







Broadband wave field evolution



$$k_p \eta$$
 -0.15 0.00

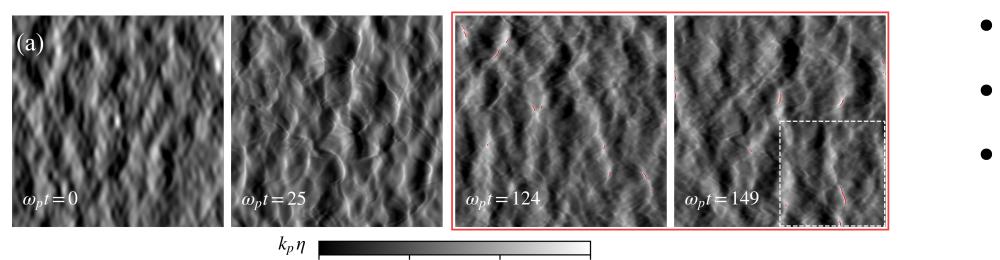
Two global steepness (slope) parameter:

- Effective slope (energy): $k_p H_s = 4k_p \langle \eta^2 \rangle^{1/2} = 4k_p (\int_0^{1/2} k_p (\int_0^{1/2} k_p (\eta^2)^{1/2} k_p (\eta^2)^{1/2})$
- Root mean square slope (roughness): $\sigma = \langle m_{\chi} \rangle^{1/2} +$

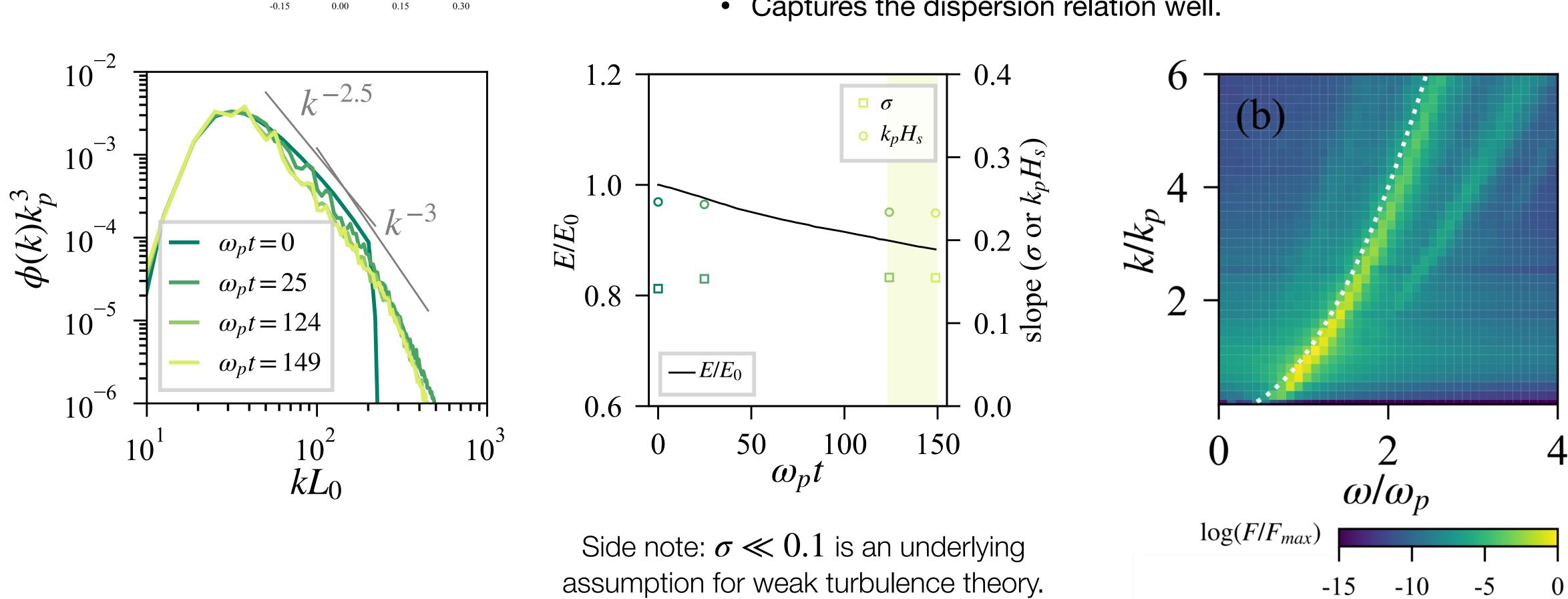
$$\int_{0}^{k_{max}} \phi(k) dk \,^{1/2}$$

- $\langle m_{y} \rangle^{1/2} = (\int_{0}^{k_{max}} k^{2} \phi(k) dk)^{1/2}, m_{x} = \partial \eta / \partial x$

Broadband wave field evolution

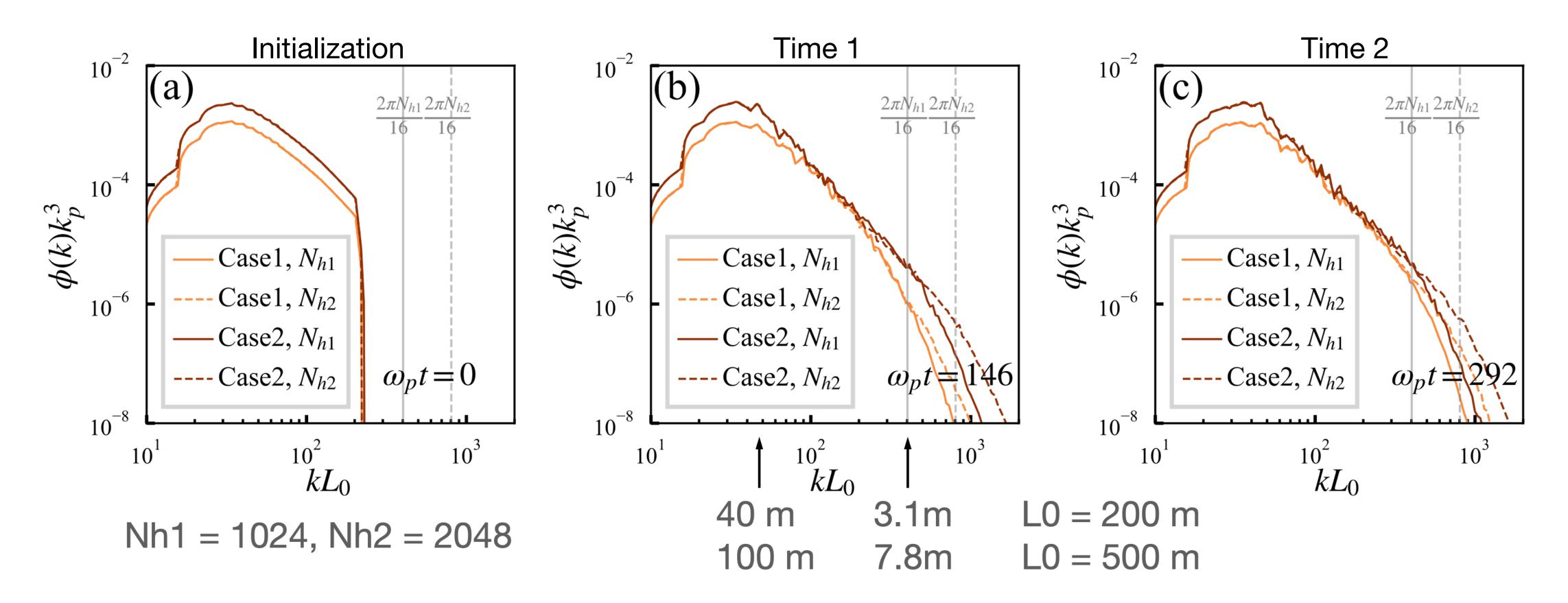


-0.15 0.00 0.15 0.30



- Saturation: k^{-3} spectral shape.
 - Energy transfer to small scale by breaking.
 - Constant energy dissipation, but the spectrum is stationary and the steepness parameters almost constant.
- Captures the dispersion relation well.

Resolution and length scales

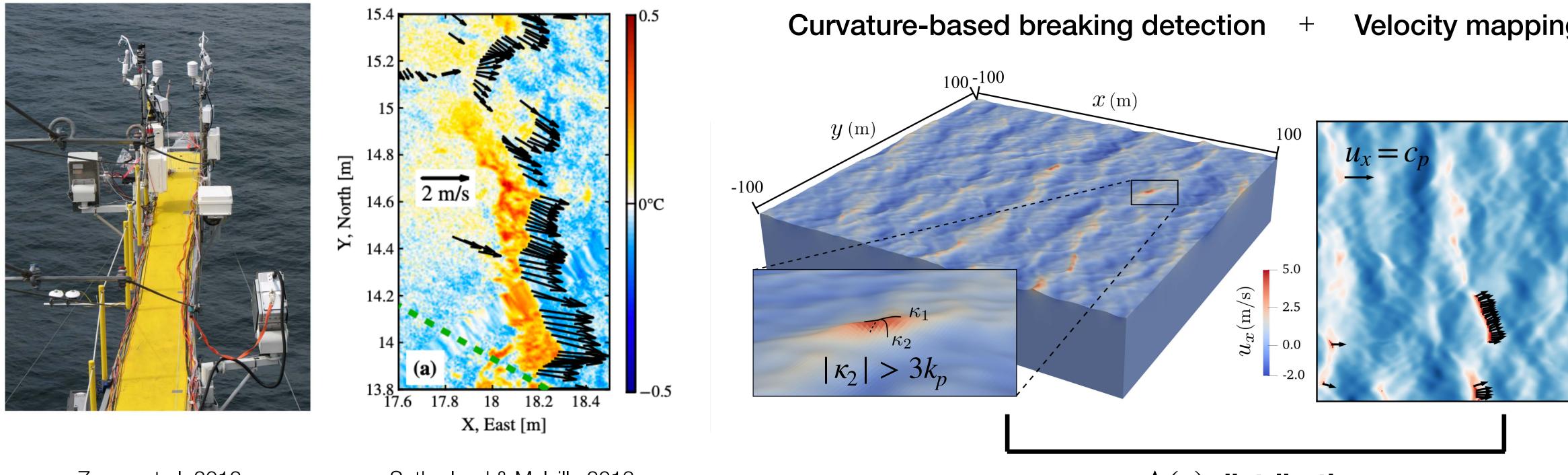


- Less steep wave field has slower transfer of energy into small scales.
- Eventually k^{-3} shape is established for numerically resolved range.

• Smallest scale is set by numerical resolution. Ratio of peak to smallest length scale O(10).

Postprocessing for $\Lambda(c)$ distribution

Field observations



Zappa et al. 2012

Sutherland & Melville 2013

Numerical experiments

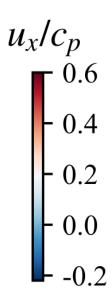
Velocity mapping

 $\Lambda(c)$ distribution

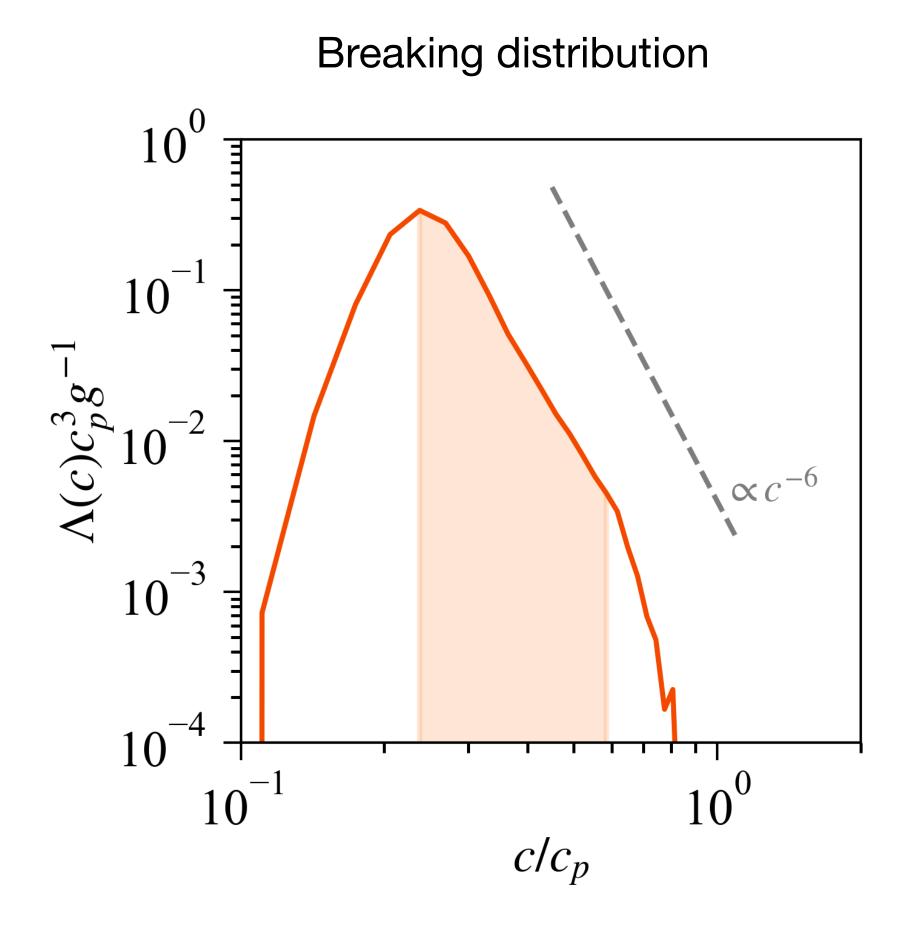
- $\Lambda(c)$: the expected length of breaking fronts (per unit area) with speed (c, c + dc).
- The variable c indicates the scale of the wave, and quantifies the strength of the breakers (affects dissipation etc.).



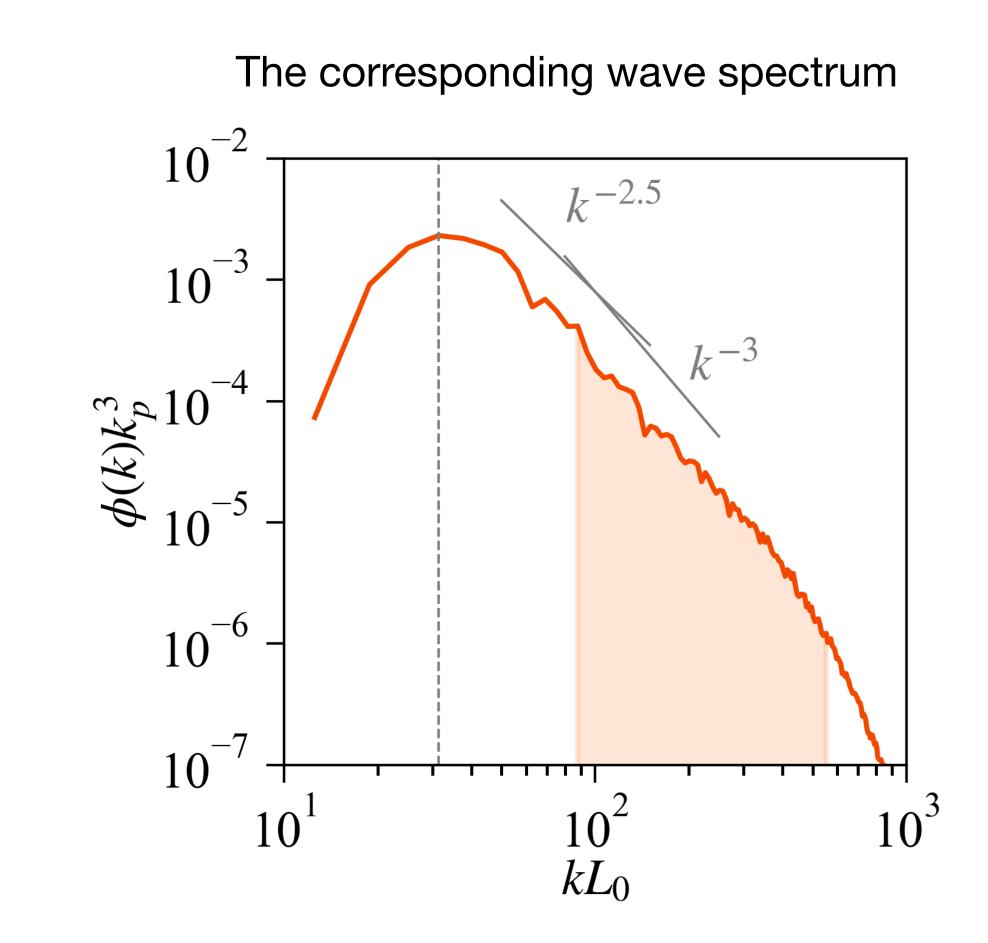




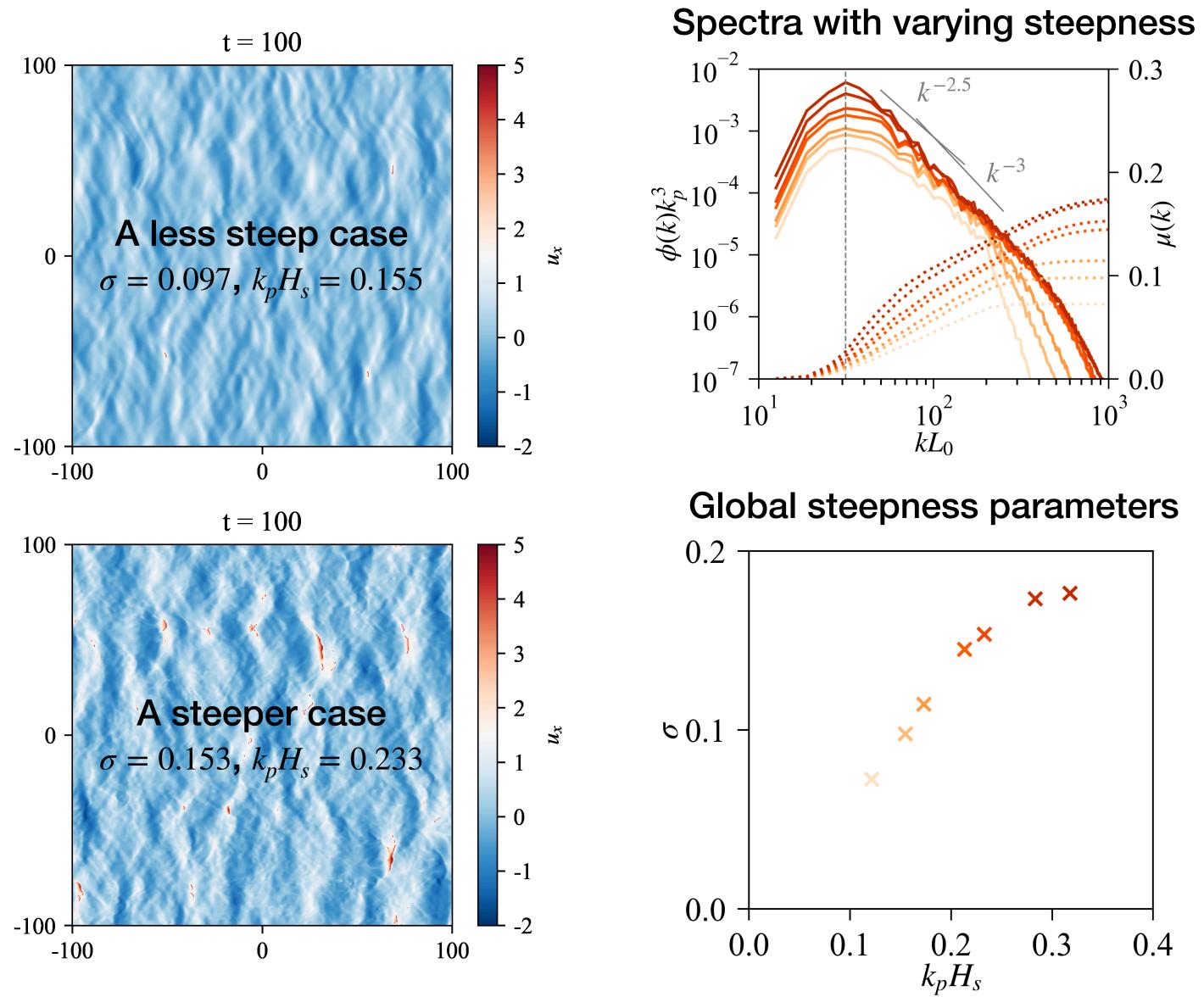
Breaking statistics - $\Lambda(c) \propto c^{-6}$ power law

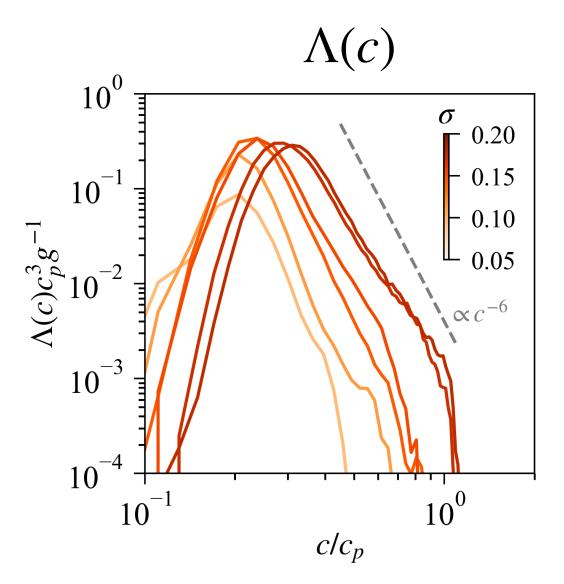


Observe $\Lambda(c) \propto c^{-6}$ power-law, which agrees with Phillips' theoretical prediction. Identify the scale of the breakers ($k = g/c^2$).

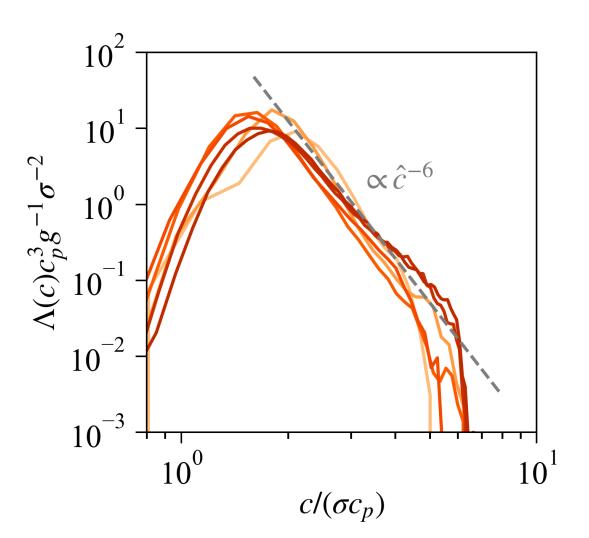


Breaking statistics - steepness scaling



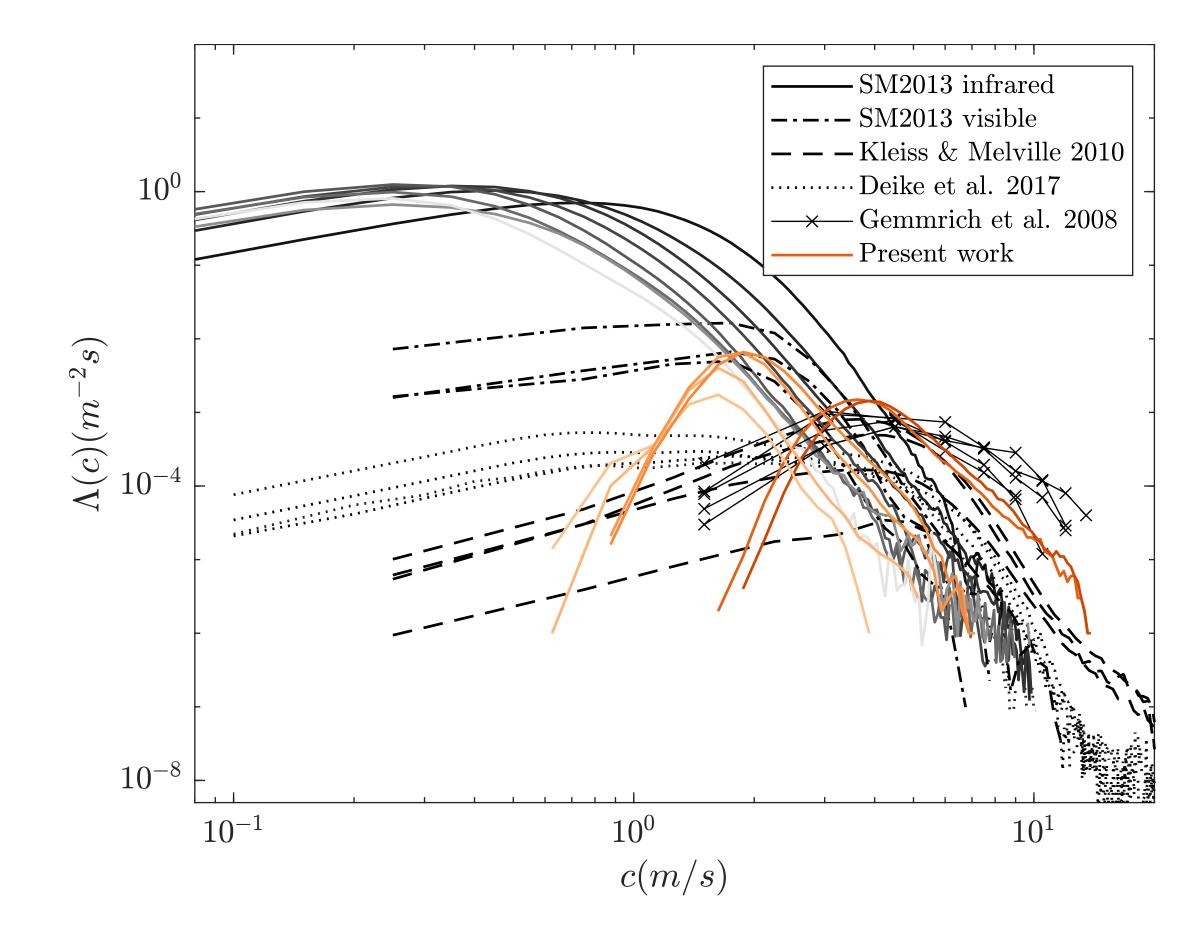


$\Lambda(c)$ with normalization



Comparison with field data

Unscaled data in physical units



Different scalings in literature

Phillips 1985 wind-only scaling (does not match data well):

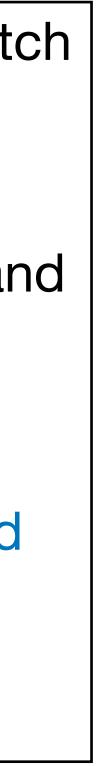
 $\Lambda(c) \propto u_*^3 g c^{-6}$

Empirical wind-wave-mixed scaling (Sutherland & Melville 2013 etc):

 $\Lambda(c)c_p^3 g^{-1} (c_p/u_*)^{1/2} \propto (c/\sqrt{gH_s})^{-6}$

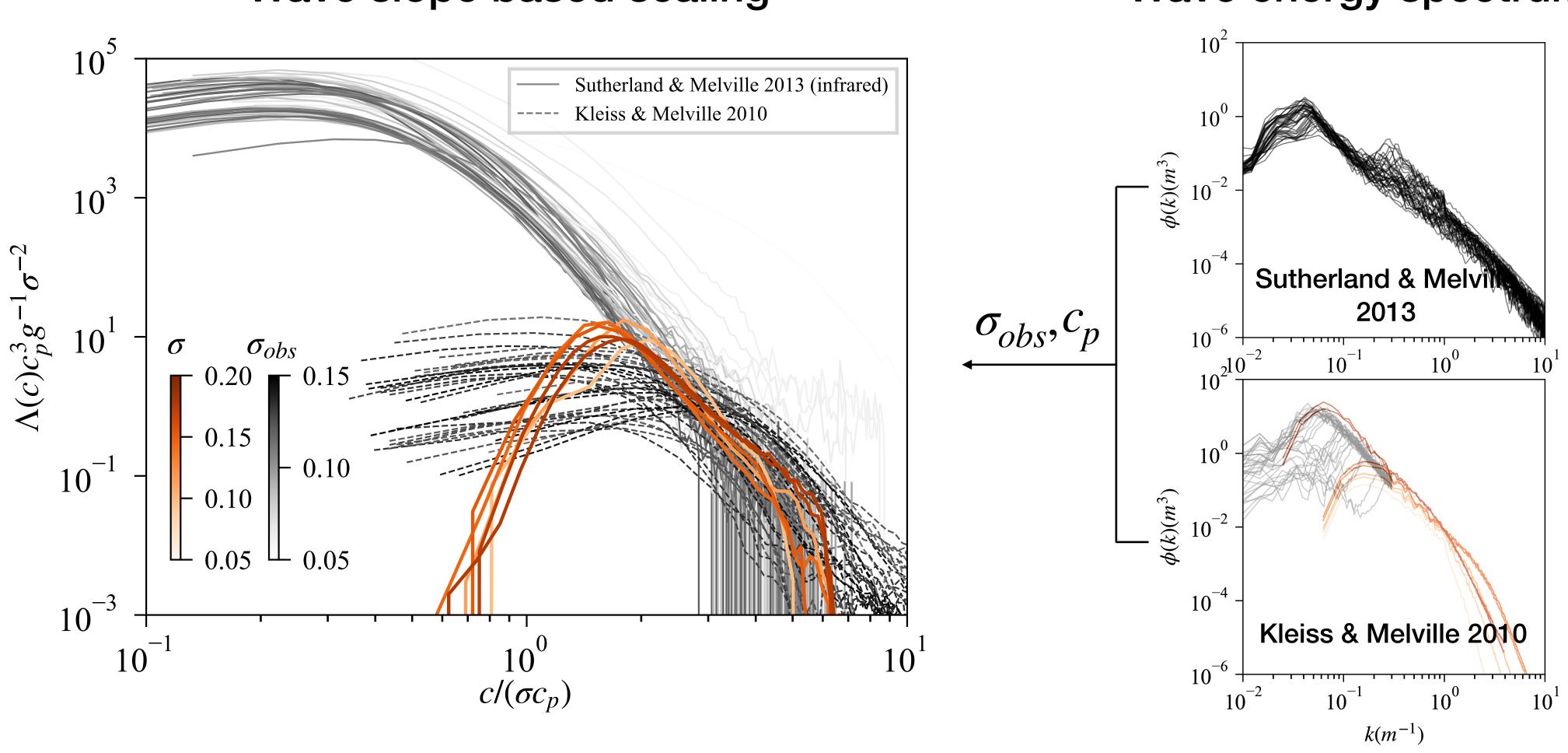
3. Present work wave-only scaling (without wind forcing):

$$\Lambda(c)c_p^3 g^{-1} \sigma^2 \propto (c/(\sigma c_p))^{-6}$$



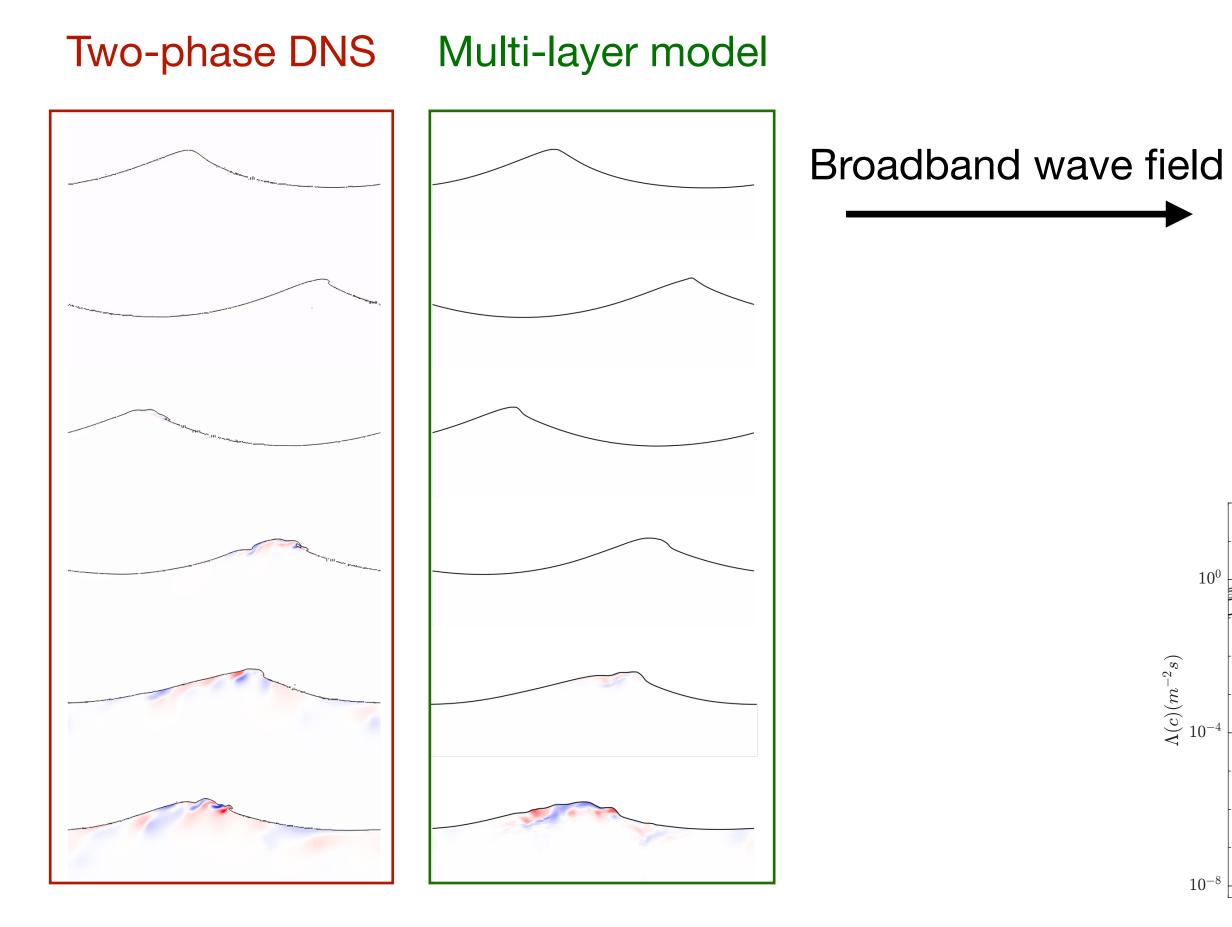
Scaling tested on field data

Wave slope based scaling



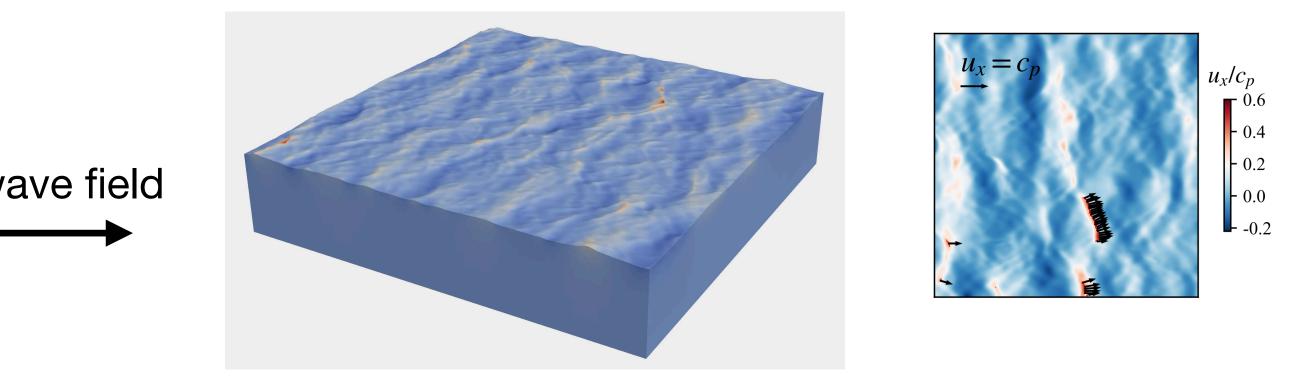
Preprint: https://arxiv.org/abs/2211.03238

Wave energy spectrum

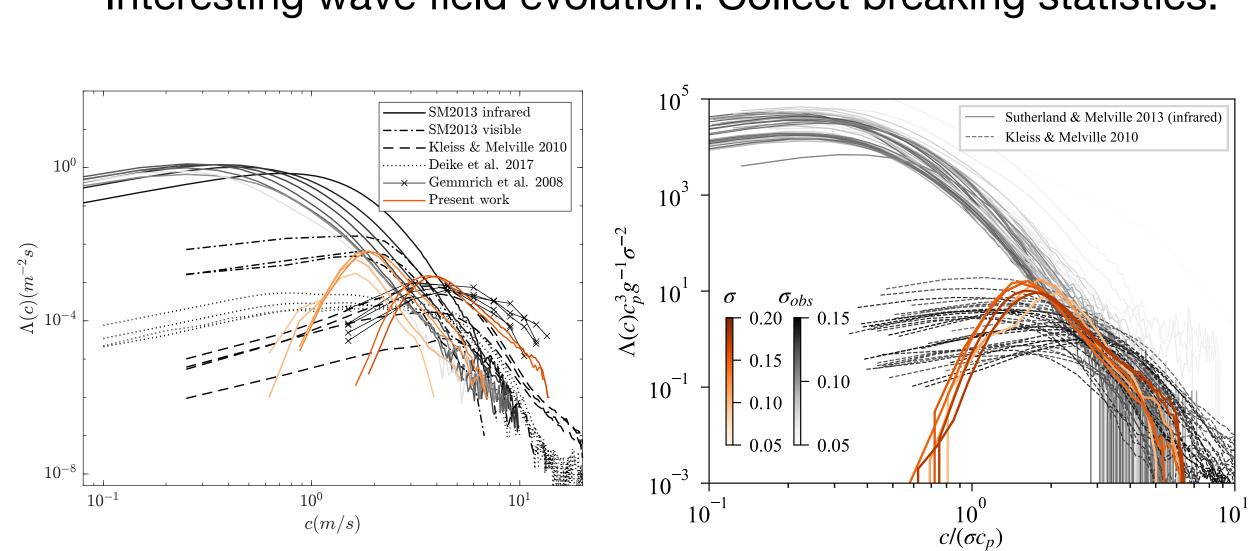


Phase-resolved wave simulations in the physical space that permit strong non-linearity (without surface overturning) with reasonable computational cost.

Summary



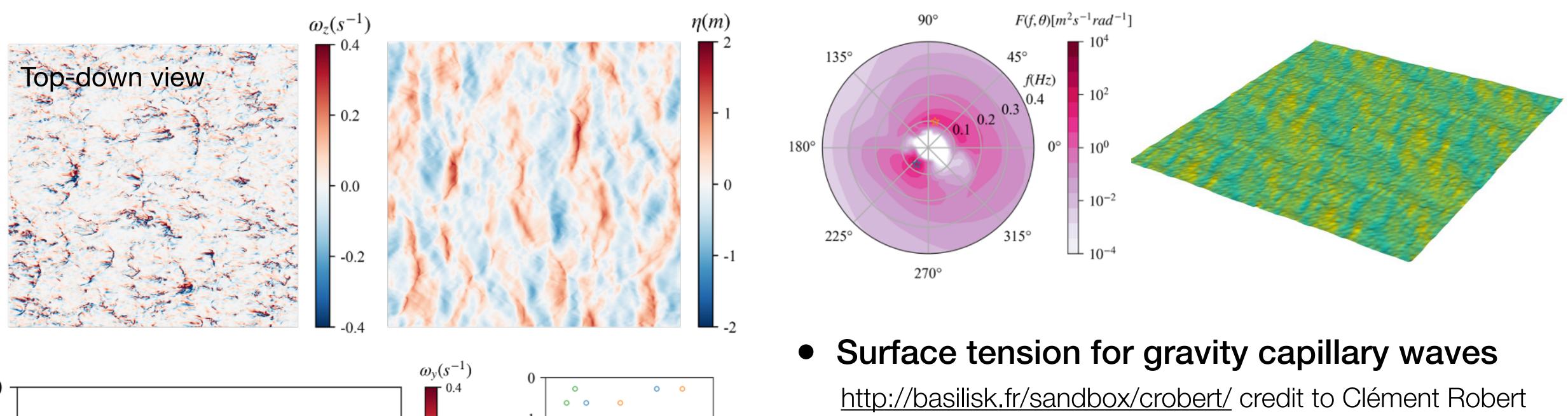
Interesting wave field evolution. Collect breaking statistics.

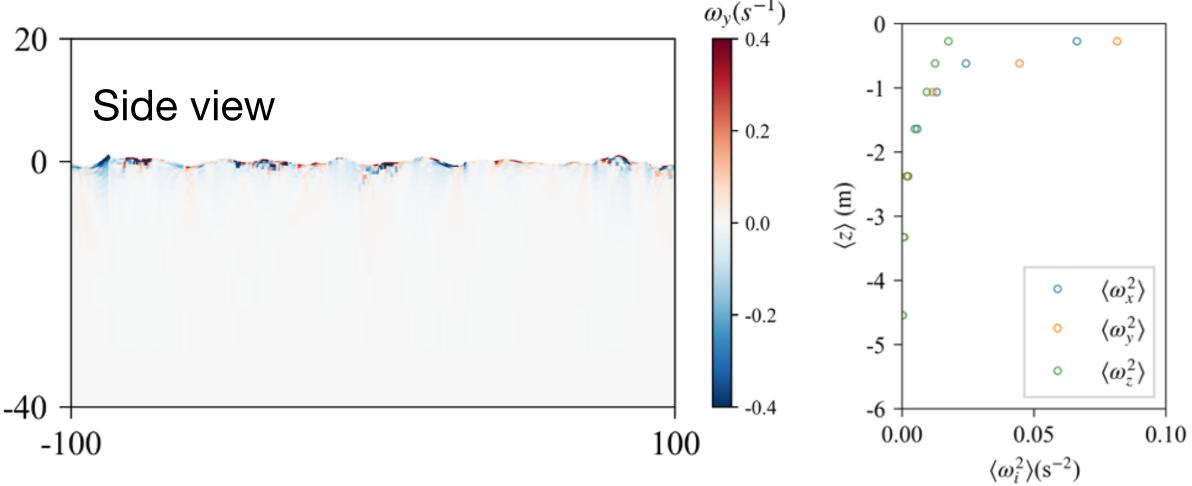


Similar level of breaking frequency to field observations. Proposed wave-slope only scaling.

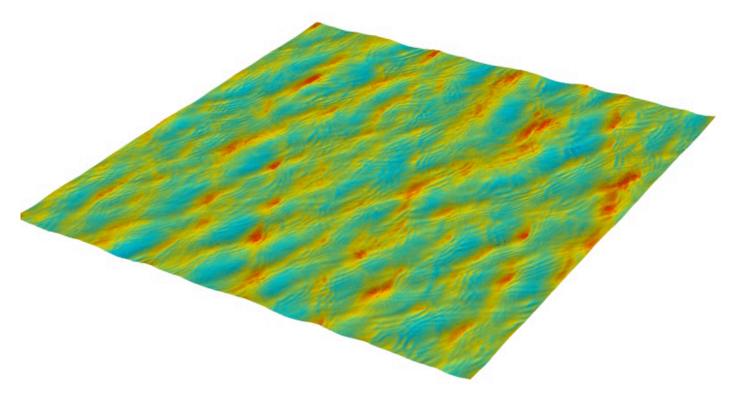
Future work

Underwater vorticity and dissipation

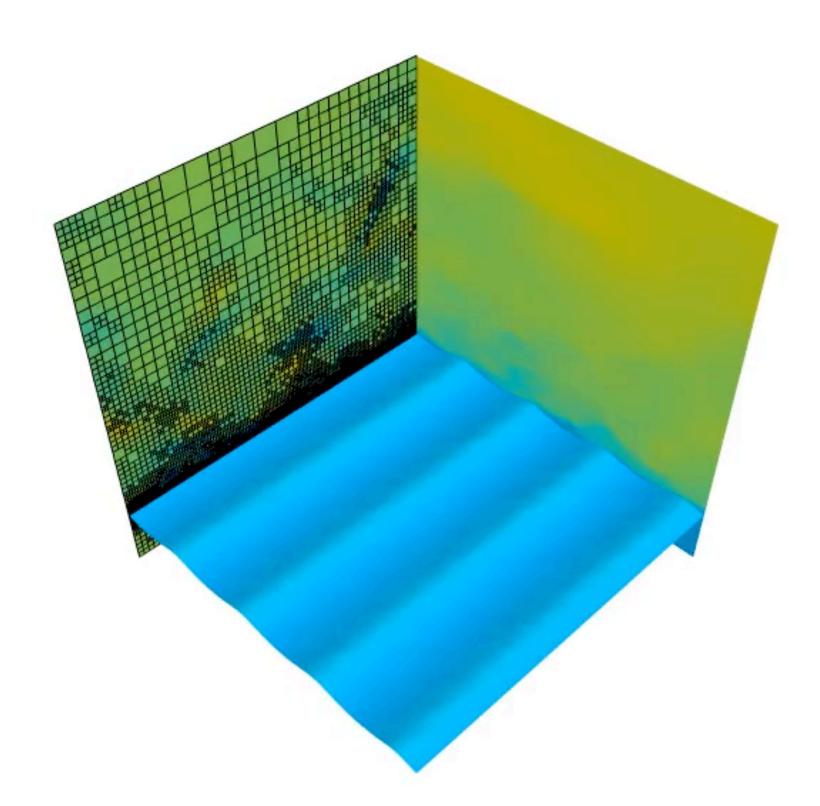


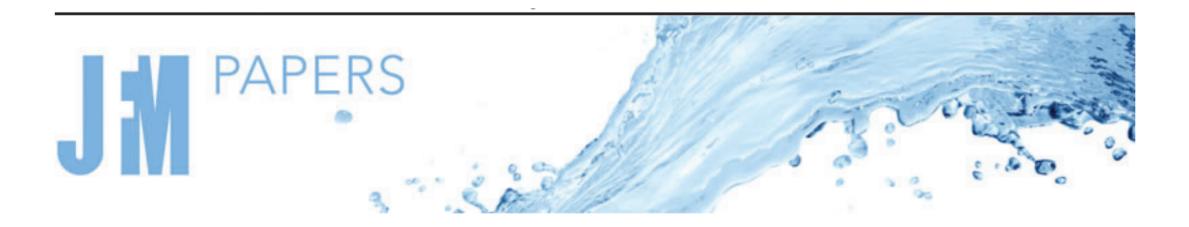


• Mixed swell-wind-wave sea state



Wind wave growth with VOF





Revisiting wind wave growth with fully coupled direct numerical simulations

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