

# High-speed impact across scales: a hybrid approach

Basilisk (and Gerris) Users' Meeting

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# Introduction

## Group updates

High-speed impact  
across scales  
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Dr. Michael Negus (2022)



Dr. Ben Fudge (2023)



Oscar Holroyd (2025)



Sebastian Dooley (2026)

Some activities in the wider collaborative group:

### Drop dynamics

- Bouncing
- Coalescence
- Splashing
- Fluid-structure interaction



### Liquid films

- Multi-physics modelling
- Asymptotic analysis
- Control theory
- Equation discovery

### Industrial mathematics

### Sustainable software

### Outreach and art



Acknowledgements:

- EPSRC (EP/V051385/1 on liquid film control, EP/W032201/1 on ReproHacks) and NSF (EP/W016036/1)
- UK Fluids Network (Drop Dynamics and Interfacial Flows SIGs)



@rcimpeanu

rcsc-group 

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- Pre-impact
- Early impact
- Post-impact

### Helmholtz flows

- Symmetric jets
- Integro-differential equation
- Inner region

### Levitating cylinders

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# Throwback Thursday



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# Motivation

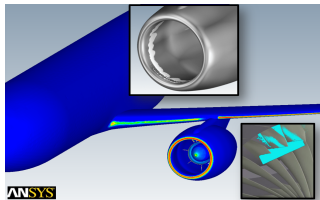
## Background

Studying the dynamics of droplets in high speed flow conditions has ramifications particularly in the aeronautics industry, where the following flight conditions:

- ▶ heavy rainfall
- ▶ high liquid water content (LWC) regions - clouds

are commonplace.

Ice formation severely affects aerodynamic performance. Of particular interest are the adverse effects on the nacelle system.





# Single drop impact

Droplet behaviour - summary - RC & Papageorgiou, IJMF 107, 192-2017 (2018)



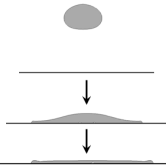
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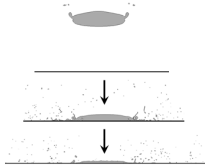
Detailed analysis of splashing dynamics has been performed in both two and three dimensions, revealing similar qualitative and quantitative features.

- ▶ small drops: pancaking behaviour, no splashing;
- ▶ medium drops: deformation before impact, moderate splashing;
- ▶ large drops: stronger deformation, violent splashing;

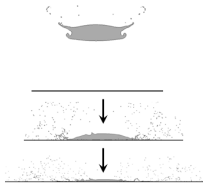
$d^* = 20 \mu\text{m}$



$d^* = 200 \mu\text{m}$

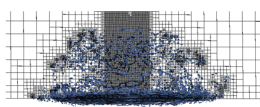
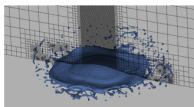
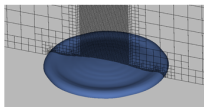


$d^* = 2000 \mu\text{m}$



2D

3D



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# Pre-impact dynamics

Experiments (NASA & INTA)

During 2012 – 2015, Vargas, Sor and Magarino (NASA - INTA Madrid collaboration) have conducted a set of experiments on droplet breakup near the leading edge of an airfoil.

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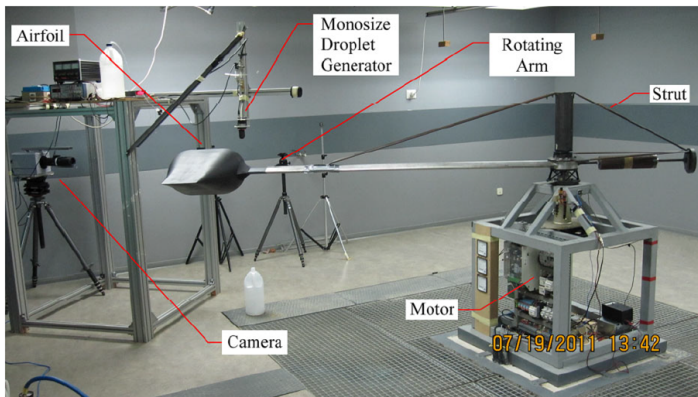


Figure 2.—Experiment set-up in the INTA test cell.

End



# Pre-impact dynamics

Experiments (NASA & INTA)

During 2012 – 2015, Vargas, Sor and Magarino (NASA - INTA Madrid collaboration) have conducted a set of experiments on droplet breakup near the leading edge of an airfoil.

Key observations:

- ▶ small drops tend to retain their shape;
- ▶ medium sized drops flatten into squashed ellipsoidal shapes;
- ▶ large drops eventually break under violent rupture.



Fig. A1 Case 1249 droplet deformation evolution: airfoil velocity = 90 m/s, chord size = 690 mm, and droplet radius = 181  $\mu\text{m}$ .

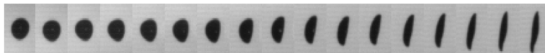
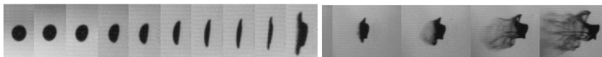


Fig. A5 Case 1041 droplet deformation evolution: airfoil velocity = 90 m/s, chord size = 690 mm, and droplet radius = 565  $\mu\text{m}$ .



Droplet deformation and breakup as  $D = 362 \mu\text{m}$  (above),  
 $D = 1130 \mu\text{m}$  (middle) and  $D = 2064 \mu\text{m}$  (below)

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# Pre-impact dynamics

Comparison - RC & Papageorgiou, IJMF 107, 192-207 (2018)

Excellent qualitative agreement is found between the experiments and computations in the same parameter range, with good quantitative agreement also found in the proposed deformation rate metric.

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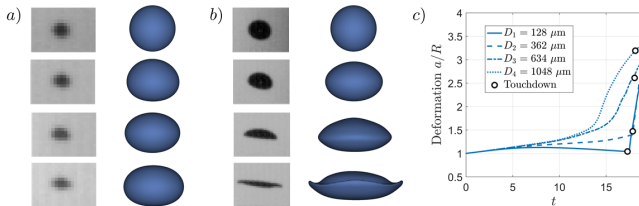
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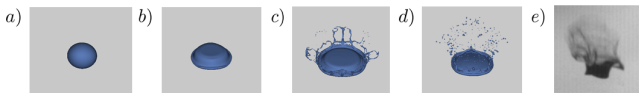
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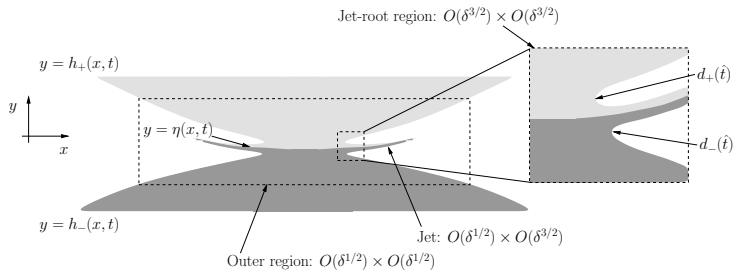
With the aid of the simulations dynamics for **much smaller drops** is analysed, while also gathering break-up data in the timesteps just before impact.



# Impact

Asymptotic structure - RC & M.R. Moore, JFM 856, 764-796 (2018)

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Early-time asymptotic structure according to Wagner theory.

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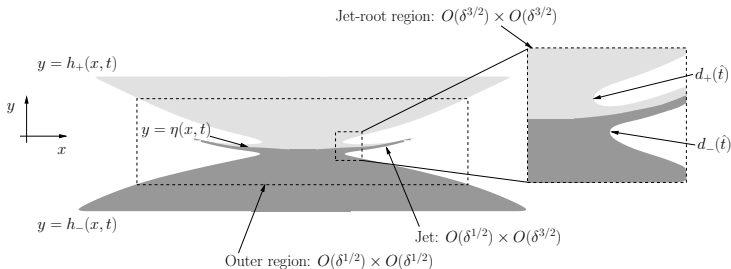
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# Impact

Asymptotic structure - RC & M.R. Moore, JFM 856, 764-796 (2018)



Early-time asymptotic structure according to Wagner theory.

- ▶ **Outer:** boundary conditions linearise onto  $y = 0$ , solved using Riemann-Hilbert techniques.
- ▶ **Inner ('jet-root'):** quasi-steady Helmholtz flow, solved using Schwarz-Christoffel mappings.
- ▶ **Jet:** thin, high-speed jet governed by zero-gravity shallow-water equations, solved using characteristics.

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# Results

Summary - RC & M.R. Moore, JFM 856, 764-796 (2018)

Given time and patience, **useful quantitative information** about the impact process and properties of the resulting jet is retrieved:

Location of jet root:

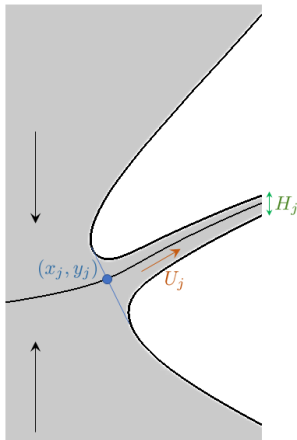
$$\hat{x}_j = 2\sqrt{\frac{(1+V)R}{1+R}}\hat{t},$$
$$\hat{y}_j = \left(\frac{(3RV+R-V-3)}{2(1+R)}\hat{t}\right).$$

Jet thickness:

$$H_j(\hat{t}) = \frac{\pi(1+V)^{3/2}}{8}\sqrt{\frac{1+R}{R}}\hat{t}^{3/2},$$

Jet velocity:

$$U_j = 2\hat{d}_0 = 2\sqrt{\frac{(1+V)R}{1+R}}\hat{t}^{-1/2},$$



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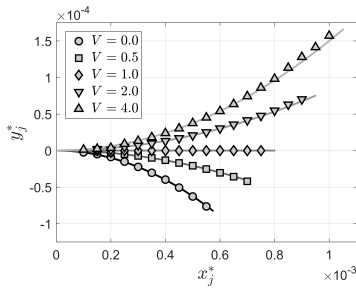
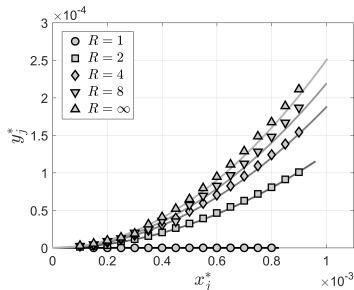
Acknowledgements

# Results

Jet root location - RC & M.R. Moore, JFM 856, 764-796 (2018)

At early times we find excellent agreement between the two approaches. This begins to deteriorate (in an anticipated manner):

- ▶ at the tip of the jet;
- ▶ when not correcting for the presence of **entrapped air bubbles**;
- ▶ once we force the underlying assumptions (e.g. lower impact velocities);
- ▶ at sufficiently large times;



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# Results

Jet angle - RC & M.R. Moore, JFM 856, 764-796 (2018)

From the inner region we can also derive the jet angle, found to be

$$\alpha = \delta^{1/2} \left( (R-1) \sqrt{\frac{1+V}{R(1+R)}} \right) \sqrt{\hat{t}} + o(\delta^{1/2}),$$

where  $\alpha$  is small. Abstractly it represents the slope of the vortex sheet in the outer region as it approaches the turnover point.  $R = 1$ ,  $V = 2$ :

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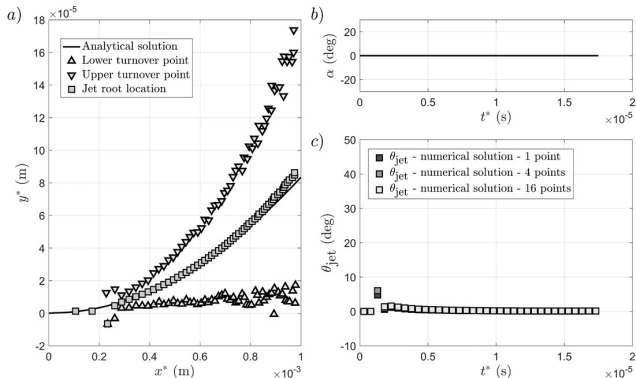
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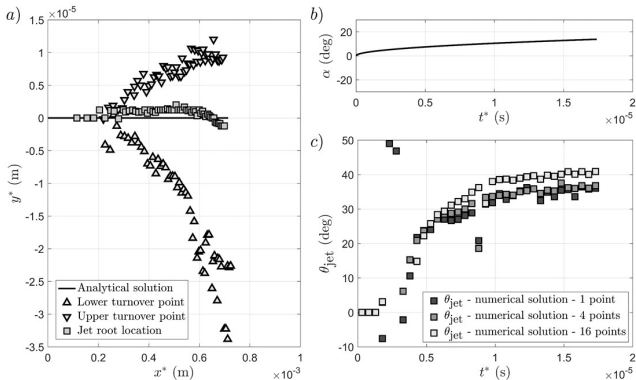
# Results

Jet angle - RC & M.R. Moore, JFM 856, 764-796 (2018)

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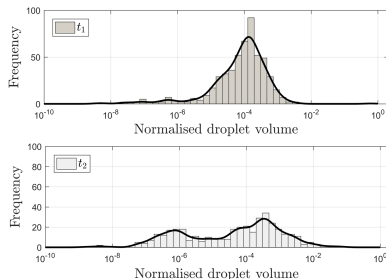
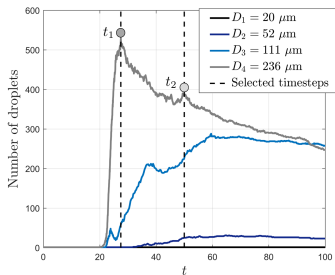
End



# Post-impact dynamics

Droplet distribution - RC & Papageorgiou, *IJMF* 107, 192-2017 (2018)

The volumes of the satellite droplets in the system follow a log-normal distribution during most of the simulation, with an average area equal to roughly 1/10000 relative to the initial droplet. This behaviour is consistent with experimental observations (Mundo, *IJMF*, 95 and Yarin, *JFM*, 1995) for lower velocity impacts onto solid surfaces.



Droplet distribution for an  $\theta = 60^\circ$  angle of incidence impact in the case of a range of droplet sizes.

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# Post-impact dynamics

Twin peaks - RC & Papageorgiou, IJMF 107, 192-2017 (2018)

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### Helmholtz flows

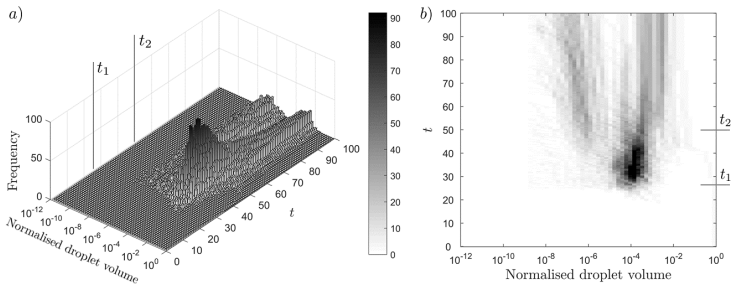
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The evolution of the drop size distribution hints at a separation in drop behaviour depending on size:

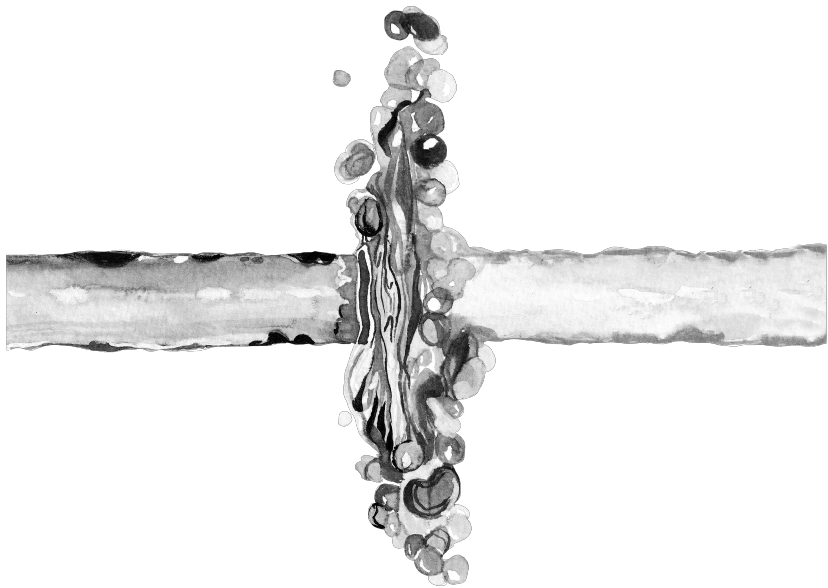
- ▶ the larger fluid volumes are found as spherical caps in contact with the solid surface.
- ▶ are more likely to grow in time as coalescence events take place.
- ▶ the smaller drops are airborne and subjected to the strong background flow.
- ▶ tend to break up as long as capillary forces allow it.

End

So . . . what have we been up to since then?



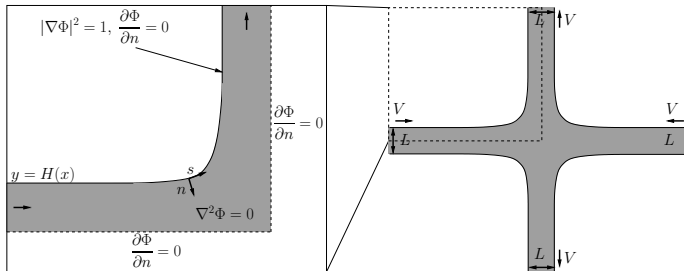
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# Results

Symmetric jets - M.R. Moore, RC, Ock<sup>2</sup>, J.M. Oliver, JFM 882, A19 (2020)

The simplest problem where we can study the effect of viscous corrections near the free surface is a pair of jets colliding symmetrically at the origin.



We find that the perturbation to the free surface satisfies the singular integro-differential equation

$$0 = \frac{d\tilde{h}_0}{ds} + 2\kappa(s) + \frac{1}{2} \int_{-\infty}^{\infty} \operatorname{cosech} \left( \frac{\pi(s-t)}{2} \right) \left[ \tilde{h}_0(t)\kappa(t) - 2\kappa(t)\tilde{w}^*(t) \right] dt,$$

where  $\tilde{w}^* = -2\arctan \left( e^{\pi s/2} \right)$ .

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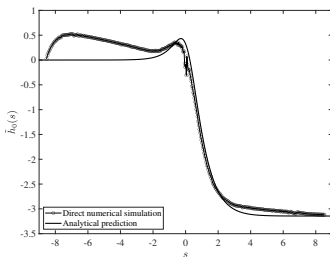
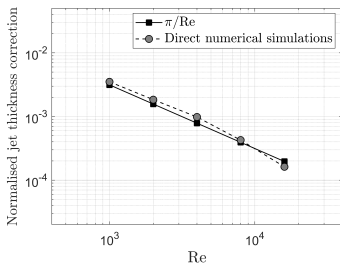
# Results

Symmetric jets - M.R. Moore, RC, Ock<sup>2</sup>, J.M. Oliver, JFM 882, A19 (2020)

The effect of viscosity is to thicken the jet at the outlet by twice the angle the tangent to the free surface turns as  $s$  increases.

- ▶ the ratio of the jet thickness at the outlet ( $s = \infty$ ) to that at the inlet ( $s = -\infty$ ) is given by

$$\frac{(H + \varepsilon^2 \tilde{h}_0)(\infty)}{(H + \varepsilon^2 \tilde{h}_0)(-\infty)} = 1 + \frac{\pi}{Re} + o\left(\frac{1}{Re}\right).$$



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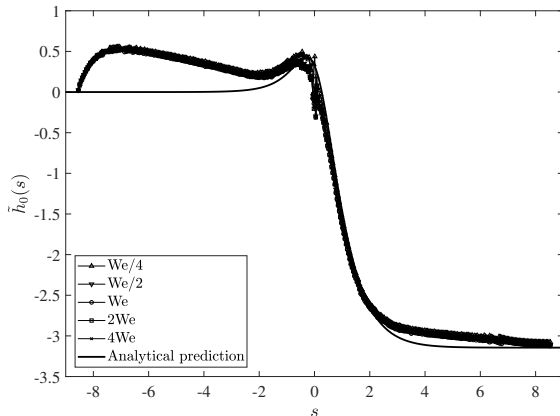
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# Results

Symmetric jets - M.R. Moore, RC, Ock<sup>2</sup>, J.M. Oliver, JFM 882, A19 (2020)

- ▶ changes in surface tension coefficient (while  $Re$  is large but finite with  $1/We = \mathcal{O}(1/Re)$ ) induce no qualitative differences;
- ▶ quantitatively, the free surface suffers a normal translation given by  $\phi_1 = -\kappa X(y - 1/We)$ .



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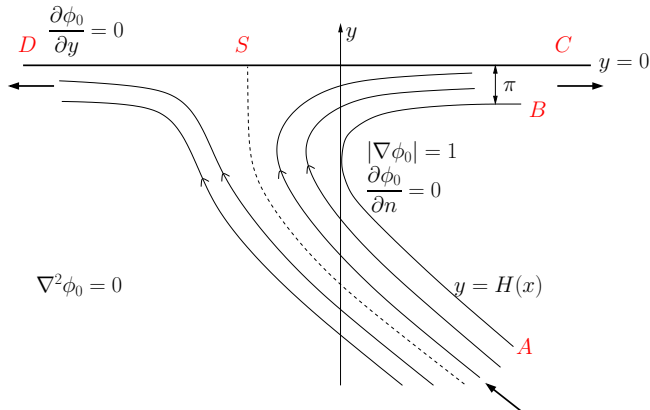
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# Results

Inner region - M.R. Moore, RC, Ock<sup>2</sup>, J.M. Oliver, JFM 882, A19 (2020)

The key example is the (suitably scaled) leading-order-inner problem according to Wagner theory for a droplet-droplet impact at small times.

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The dashed line indicates a dividing streamline that terminates at a relative stagnation point on the body at S. In a frame moving with the turnover points, the problem is a Helmholtz flow.

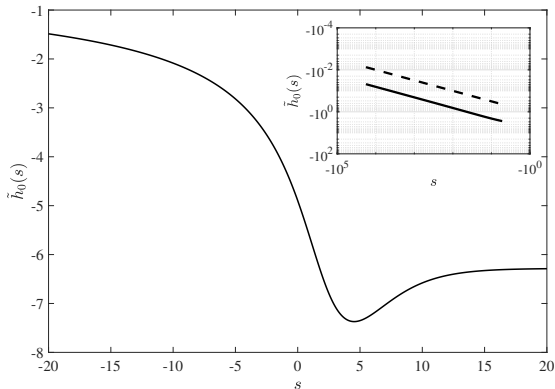
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# Results

Inner region - M.R. Moore, RC, Ock<sup>2</sup>, J.M. Oliver, JFM 882, A19 (2020)

- ▶ Downstream as  $s \rightarrow \infty$ , one can see the perturbation approaching an analytically predicted value of  $-2\pi$ ;
- ▶ Upstream we see inverse square-root decay back into the bulk (the outer Wagner region), displayed in further detail in the inset.



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Inner region - M.R. Moore, RC, Ock<sup>2</sup>, J.M. Oliver, JFM 882, A19 (2020)

Despite the problem being in 2D, numerically this is a challenge:

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Inner region - M.R. Moore, RC, Ock<sup>2</sup>, J.M. Oliver, JFM 882, A19 (2020)

Despite the problem being in 2D, numerically this is a challenge:

- ▶ initial and boundary conditions are implemented based on asymptotic expressions valid 'in the far-field';



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Inner region - M.R. Moore, RC, Ock<sup>2</sup>, J.M. Oliver, JFM 882, A19 (2020)

Despite the problem being in 2D, numerically this is a challenge:

- ▶ initial and boundary conditions are implemented based on asymptotic expressions valid 'in the far-field';
- ▶ the underlying theory is powerful and valid up to second order



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- ▶ initial and boundary conditions are implemented based on asymptotic expressions valid 'in the far-field';
- ▶ the underlying theory is powerful and valid up to second order, but ...



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Despite the problem being in 2D, numerically this is a challenge:

- ▶ initial and boundary conditions are implemented based on asymptotic expressions valid 'in the far-field';
- ▶ the underlying theory is powerful and valid up to second order, but ... it ignores the gas altogether, and it need not bother with contact angles, outflow conditions etc.



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Inner region - M.R. Moore, RC, Ock<sup>2</sup>, J.M. Oliver, JFM 882, A19 (2020)

Despite the problem being in 2D, numerically this is a challenge:

- ▶ initial and boundary conditions are implemented based on asymptotic expressions valid 'in the far-field';
- ▶ the underlying theory is powerful and valid up to second order, but ... it ignores the gas altogether, and it need not bother with contact angles, outflow conditions etc.
- ▶ the higher the  $Re$ , the better the analytical predictions.



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# Results

Inner region - M.R. Moore, RC, Ock<sup>2</sup>, J.M. Oliver, JFM 882, A19 (2020)

Despite the problem being in 2D, numerically this is a challenge:

- ▶ initial and boundary conditions are implemented based on asymptotic expressions valid 'in the far-field';
- ▶ the underlying theory is powerful and valid up to second order, but ... it ignores the gas altogether, and it need not bother with contact angles, outflow conditions etc.
- ▶ the higher the  $Re$ , the better the analytical predictions. The higher the  $Re \Rightarrow$  ill-conditioning and instabilities.



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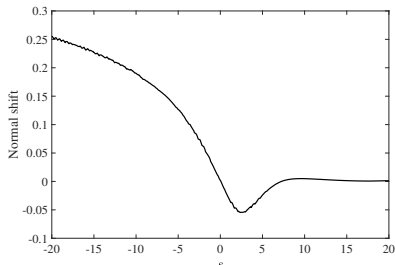
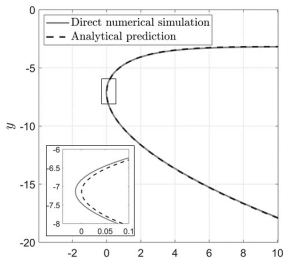
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A truly symbiotic relationship?



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# In the same spirit ...

Levitating a cylinder - Dalwadi, RC et al., JFM 917, A28 (2021)

Direct numerical simulations can be deployed alongside beautiful matched asymptotic expansion methods for a wide class of problems as means to resolve complex physical systems with no fitting parameters, with experiments providing essential confirmation.

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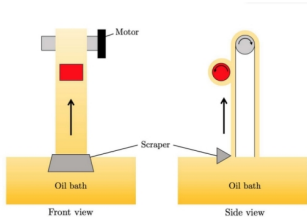
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*J. Fluid Mech.* (2021), vol. 917, A28, doi:10.1017/jfm.2021.284



## Levitation of a cylinder by a thin viscous film

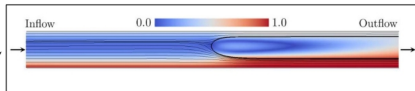
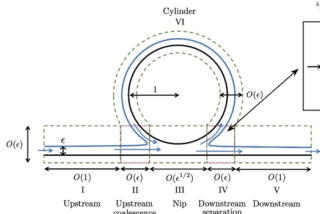
Mohit P. Dalwadi<sup>1,†</sup>, Radu Cîmpeanu<sup>1,2,3</sup>, Hilary Ockendon<sup>1</sup>,  
John Ockendon<sup>1</sup> and Tom Mullin<sup>1,4</sup>

<sup>1</sup>Mathematical Institute, University of Oxford, Oxford OX2 6GG, UK

<sup>2</sup>Mathematical Institute, University of Warwick, Zeeman Building, Coventry CV4 7AL, UK

<sup>3</sup>Department of Mathematics, Imperial College London, London SW7 2AZ, UK

<sup>4</sup>Linacre College, University of Oxford, Oxford OX1 3JA, UK



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Bringing it all together - a hybrid multi-scale framework



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# Hybrid framework

## Multi-scale methodology

For us to be able to get closer to an (accurate and efficiently obtained) answer, any solution would need to resolve:

- ▶ **sub-micron scales** on impact and for the fragmentation process.
- ▶  $\mathcal{O}(1)$  m or larger lengthscales for **device-specific geometries**.



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No monolithical solver could viably do this.



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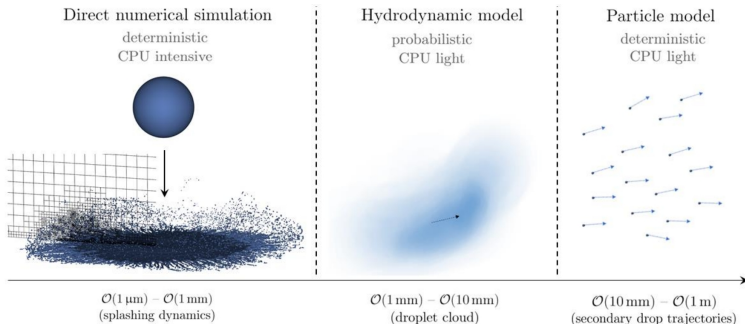
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# Acknowledgements



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# Conclusions

## Acknowledgements

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- ▶ the Oxford Centre for Industrial and Applied Mathematics,
- ▶ the Warwick Mathematics Institute,

without which this work would not have been possible.



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